

## Raman study on structure of $U_{1-y}Gd_yO_{2-x}$ ( $y=0.005, 0.01, 0.03, 0.05$ and $0.1$ ) solid solutions



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### HIGHLIGHTS

- Investigation of structural character of  $U_{1-y}Gd_yO_{2-x}$  solid solutions.
- Defect structures in  $U_{1-y}Gd_yO_{2-x}$  solid solutions were evaluated by Raman spectroscopy.
- Oxygen deficiency due to Gd content causes Raman band related to oxygen vacancy.

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### ABSTRACT

The  $U_{1-y}Gd_yO_{2-x}$  solid solutions with  $y = 0.005, 0.01, 0.03, 0.05$  and  $0.1$  were characterized by Raman spectroscopy to investigate the defect structure induced by oxygen vacancies. The oxygen deficiencies of solid solutions were estimated by the relation between the doping level and a lattice parameter calculated from X-ray diffraction patterns. Raman mode shifts to higher wavenumber with increasing doping level showed that crystal lattice disorder due to oxygen vacancies. The frequency shifts and relative ratio of Raman modes were enabled to be the indicator for composition, defect and oxygen vacancy of  $U_{1-y}Gd_yO_{2-x}$  solid solutions.

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### 1. Introduction

In reactor irradiation of  $UO_2$  nuclear fuel leads to the formation of a wide range of fission products, transuranium elements and activation products [1–3]. Among them rare earth (RE) elements form  $U_{1-y}RE_yO_2$  solid solutions with  $UO_2$  [4]. When irradiated under reducing circumstances, to a high burnup in the range of 70–80 GWd/tU, the  $UO_2$  fuel becomes doped with fission products

(FPs), especially rare earths, and can be considered as a slightly sub-stoichiometric ( $U_{1-y}FP_yO_{2-x}$ ) or stoichiometric compound [5]. A knowledge of the structural character of RE-doped  $UO_{2\pm x}$  is very important not only to understand the characteristics of spent nuclear fuel but also to describe the thermodynamic properties and the phase relations in U-RE-O systems [6–9]. Among the many RE elements, Gd is one of the major fission products formed in solid solution with  $UO_2$  and has often been selected as a dopant in simulated spent nuclear fuels [10–12]. It is also used as a burnable absorber for the  $UO_2$ - $Gd_2O_3$  fuels developed to extend the length of the fuel cycle for PWRs [13]. Thus, Gd-doped  $UO_2$  has been extensively researched using several experimental techniques, generally in tandem with X-ray diffraction (XRD) [14–20], to understand its physical and chemical properties.

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Recently, Raman spectroscopy has been widely applied to characterize nuclear fuel materials, because it is a convenient, sensitive and nondestructive method [21]. For example, the effect of oxygen stoichiometry on the defect structure of  $\text{UO}_2$  [22,23], the oxygen sublattice structure in thorium dioxide-uranium dioxide fuel materials [24] and the influence of trivalent-dopants on the structural properties of  $\text{UO}_2$  [25] have been investigated in detail using Raman spectroscopy. The role of oxygen vacancies on the kinetics of oxidation of fuel have also been demonstrated in Raman studies of Gd-, Dy-doped  $\text{UO}_2$  [25],  $\text{U}_{1-y}\text{Nd}_y\text{O}_{2-x}$  [26],  $\text{U}_{1-y}\text{La}_y\text{O}_{2-y/2}$  [27], and  $\text{U}_{1-y}\text{Am}_y\text{O}_{2-x}$  [28] samples. The oxygen vacancy can be created to maintain electroneutrality in RE(III)-doped  $\text{UO}_2$  when RE(III) is substituted for U(IV) in solid solutions.

In this study we have characterized the effect of Gd-doping on the structure ( $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$ ) using XRD and Raman spectroscopy to identify the possible defect structures in nuclear fuel materials. Defect structures of  $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$  solid solutions with various Gd doping levels were analyzed, and the results compared with published literature.

## 2. Experimental

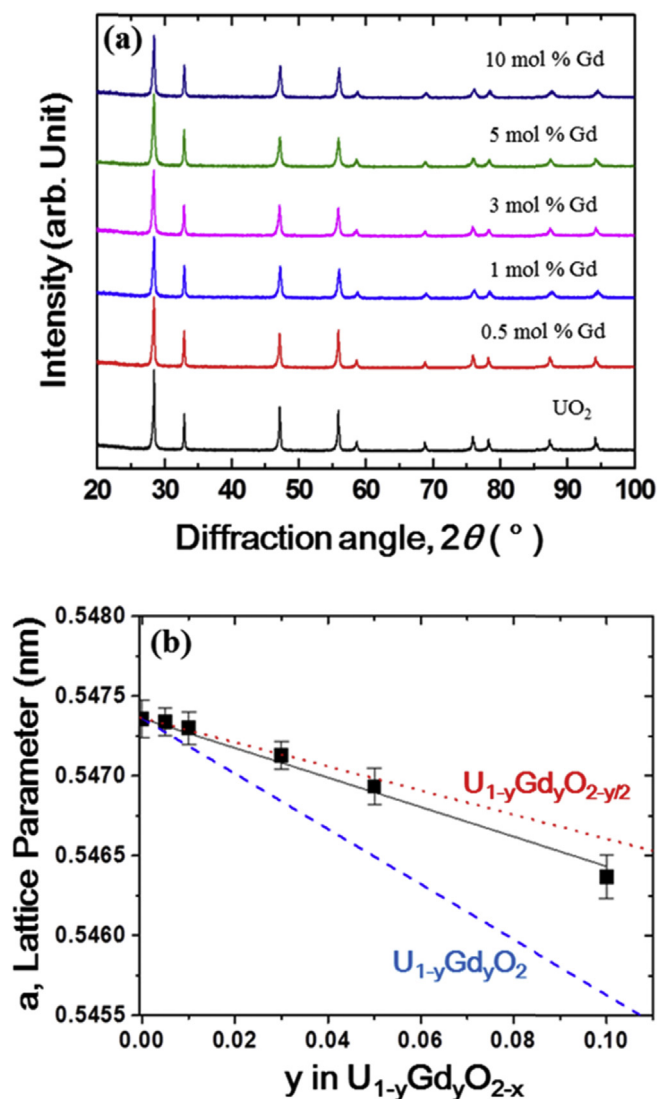
$\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$  solid solution pellets with various composition ( $y = 0.005, 0.01, 0.03, 0.05$  and  $0.1$ ) were synthesized by a conventional solid-state reaction with powder mixing. Appropriate amounts of  $\text{UO}_2$  and  $\text{Gd}_2\text{O}_3$  (Aldrich, >99.99%) powders to achieve the intended composition were blended thoroughly in an agar mortar. The mixtures were then pressed into a disk-shaped pellet and sintered in an alumina tube furnace (Ajeon Heating Industrial, Korea) at  $1700^\circ\text{C}$  for 18 h under a reducing atmosphere with flowing  $\text{H}_2$  to produce  $\text{U}^{4+}_{1-y}\text{Gd}^{3+y}\text{O}^{2-}_{2-y/2}$  type solid solutions [8,15]. The sintered pellets were cooled to room temperature in flowing  $\text{H}_2$  after annealing in the same atmosphere at  $1200^\circ\text{C}$  for 12 h. An undoped  $\text{UO}_2$  pellet was also prepared using the same procedure. After sintering, the pellets were stored in a vacuum chamber to prevent surface oxidation before measuring XRD and Raman spectra.

The X-ray diffraction (XRD) patterns of the  $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$  solid solution pellets were performed using a Bruker AXS D8 Advance X-ray Diffractometer using  $\text{CuK}\alpha$  radiation at room temperature. XRD data was collected in the range from  $20^\circ$  to  $120^\circ$  using a  $0.02^\circ$  step size. These analyses required the exposure of the specimens to air for a total time of only 500 s. Data was collected from several locations to confirm the solid structure was homogeneous. The lattice parameters were calculated by a refinement process using the TOPAS program (Bruker Analytical X-Ray Systems) with the  $Fm\bar{3}m$  space group.

The Raman spectra were measured with an ANDOR Shamrock SR500i spectrometer, with active vibrations excited using a He-Ne laser with a wavelength of 632.8 nm. The laser with *c.a.* 5 mW power was focused onto the pellets using an Olympus microscope with a 50-fold magnification lens. This laser power was confirmed as low enough to prevent surface oxidation of the pellets due to local heating by the laser beam [29]. Raman spectra were acquired over the wavenumber range  $400\text{--}1200\text{ cm}^{-1}$  at room temperature with an exposure time to air of 300 s. Raman spectra were measured at different locations on the surface of a pellet to confirm the homogeneity of the pellet and the reproducibility of the spectra.

## 3. Results and discussion

XRD patterns for  $\text{UO}_2$  and a number of  $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$  pellets with different Gd contents are shown in Fig. 1(a). All the specimens exhibit the fluorite structure and no XRD peaks for monoclinic



**Fig. 1.** (a) XRD patterns of  $\text{UO}_2$  and  $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$  solid solutions with  $y = 0.005, 0.01, 0.03, 0.05$  and  $0.1$ . (b) The lattice parameters obtained from XRD patterns of  $\text{UO}_2$  and  $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$  solid solutions with increasing Gd concentration. The red dotted and blue dashed lines show the linear relationships obtained for  $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-y/2}$  and  $\text{U}_{1-y}\text{Gd}_y\text{O}_2$  solid solutions, respectively, as the Gd content changes [15]. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

$\text{Gd}_2\text{O}_3$  and undoped  $\text{UO}_2$  were observed in the range  $25^\circ < 2\theta < 35^\circ$  in contrast to patterns recorded on  $\text{Gd}_2\text{O}_3$ -dispersed  $\text{UO}_2$  in which  $\text{Gd}_2\text{O}_3$  particles are dispersed within the  $\text{UO}_2$  matrix [30]. The lattice parameters of the pellets obtained from these patterns are plotted as a function of Gd content in Fig. 1(b). In contrast to the minor change in lattice parameter observed with increasing Gd content for the  $\text{Gd}_2\text{O}_3$ -dispersed  $\text{UO}_2$  [30], the lattice parameter for the series  $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$  solid solutions decreases linearly with the increase in Gd content. Our linear relationship for  $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$  solid solutions is less steep than that observed for  $\text{U}_{1-y}\text{Gd}_y\text{O}_2$ -type solid solutions, but well-matched with that of  $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$  ( $x \approx y/2$ )-type solid solutions [15]. Lanthanide-doped  $\text{UO}_2$  ( $\text{U}_{1-y}\text{La}_y\text{O}_{2-y/2}$ ) solid solutions also exhibited a linear relationship between the lattice parameter and the lanthanum content [27]. This analysis shows our pellets are sub-stoichiometric  $\text{U}_{1-y}\text{Gd}_y\text{O}_{2-x}$  ( $x \approx y/2$ ) solid solutions in the given range  $y = 0.005\text{--}0.05$ .

Although it is quite difficult to measure the accurate oxygen-to-metal (O/M) ratio for rare earth-doped uranium dioxides, Ohmichi

et al. [15] deduced the relationship between the lattice parameter (a) and the value of x in  $U_{1-y}Gd_yO_{2-x}$  type solid solutions with constant y to be  $da/dx = 0.024 \pm 0.006$  nm. This equation was derived using O/M ratio data by chemical analysis and the lattice parameter values by XRD analysis. Using this relationship, the O/M ratios of the sintered  $U_{1-y}Gd_yO_{2-x}$  pellets is estimated and plotted as a function of Gd content in Fig. 2. Those values closely match the results obtained for  $U_{1-y}Gd_yO_{2-x}$  ( $x \approx y/2$ )-type solid solutions in the published study [15]. The  $U_{0.9}Gd_{0.1}O_{2-x}$  pellet, however, has a higher O/M ratio than expected due possibly to oxidation during handling or a lack of reducing power during sintering.

Fig. 3 shows Raman spectra recorded on the  $UO_2$  and  $U_{1-y}Gd_yO_{2-x}$  ( $y = 0.005, 0.01, 0.03, 0.05$  and  $0.1$ ) specimens. The spectra are the averaged spectra from 10 different locations on the pellets. Some inhomogeneity was observed for spectra from different locations on the  $U_{1-y}Gd_yO_{2-x}$  pellets due to the dry mixing process that leads to some inhomogeneous distribution of the cations [16]. However, as will be shown below, these differences do not disguise the ability to observe the changes in the spectra as Gd content varies.

Raman spectra recorded on undoped  $UO_2$  exhibit the spectral features expected for an undistorted fluorite structure at  $445\text{ cm}^{-1}$  and at  $1150\text{ cm}^{-1}$ . The peak at  $445\text{ cm}^{-1}$  can be attributed to the triply degenerate Raman active  $T_{2g}$  mode for the U-O symmetric stretching mode in the fluorite structure of  $UO_2$  [31–33]. The broad peak at  $1150\text{ cm}^{-1}$  has been assigned to an overtone (2L-O) of the first order longitudinal optical (L-O) phonon regarded as a fingerprint for a quasi-perfect fluorite structure [22,34]. The very shallow feature located at  $575\text{ cm}^{-1}$  can be ascribed to a first order L-O phonon at the center of the Brillouin zone [35].

Raman spectra recorded on the  $U_{1-y}Gd_yO_{2-x}$  pellets showed the development of a broad band in the spectral region  $500$  to  $650\text{ cm}^{-1}$ , and a dramatic decrease in intensity of peaks at  $445\text{ cm}^{-1}$  and at  $1150\text{ cm}^{-1}$  as the Gd doping level increased. These spectral changes are similar to those observed for hyper-stoichiometric  $UO_{2+x}$ , due to the defects introduced by hyper-stoichiometric oxygen [23], and for La-doped  $UO_2$  due to the distortion of lattice structure induced by doping [27]. The decreasing  $T_{2g}$  and 2L-O modes with increasing Gd contents indicate distortion of the fluorite  $UO_2$  lattice structure due to the  $Gd^{3+}$  doping. The Raman spectra were analyzed by deconvoluting the broad band in the spectral region from  $500$  to  $650\text{ cm}^{-1}$  into three bands (Fig. 4). The

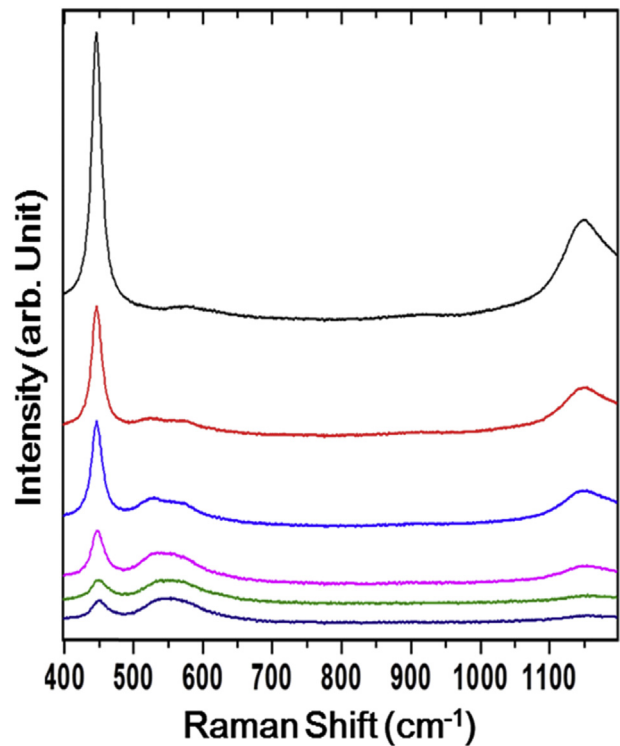


Fig. 3. Raman spectra of  $UO_2$  and  $U_{1-y}Gd_yO_{2-x}$  solid solutions with  $y = 0.005, 0.01, 0.03, 0.05$  and  $0.1$  from top to bottom.

peaks at  $\sim 445, \sim 530, \sim 575$  and  $\sim 630\text{ cm}^{-1}$  are denoted as  $T_{2g}, \alpha, \beta$  and  $\gamma$ , respectively. The frequency shifts for each peak varied slightly with increasing Gd content, as observed previously for  $(U,Am)O_{2-\delta}$  samples [28]. The averaged frequencies used to fit the broad bands in the  $500$  to  $650\text{ cm}^{-1}$  region for all measured Raman spectra are plotted as a function of Gd content for the  $T_{2g}, \alpha$  and  $\beta$  peaks in Fig. 5.

The peak  $T_{2g}$  at  $\sim 445\text{ cm}^{-1}$  in Raman spectra of  $U_{1-y}Gd_yO_{2-x}$  pellets shows a shift to higher wavenumbers is accompanied by a marked decrease in intensity with increasing Gd doping level. Raman spectra recorded on RE(III)-doped  $UO_2$  [25] and  $CeO_2$  [36] showed similar features due to the structural distortions of the fluorite lattice caused by the RE(III) doping. McBride et al. [36] suggested this frequency shift for the  $T_{2g}$  mode  $Ce_{1-y}RE_yO_{2-x}$  samples is caused by contraction of the lattice, and can be attributed to the creation of oxygen vacancies to offset the charge imbalance due to RE(III) doping. In our case, the decrease in lattice parameter (a) with increasing Gd doping indicates a contraction of the  $UO_2$  lattice and an increase in the oxygen vacancy concentration required to balance the charge of the  $U_{1-y}Gd_yO_{2-x}$  solid solution. These changes in lattice structure of the  $UO_2$  result in the positive frequency shift of the  $T_{2g}$  mode with increasing Gd content. This frequency shift may be useful in quantitative analysis of  $U_{1-y}Gd_yO_{2-x}$  solid solutions. The linear dependence of the  $T_{2g}$  mode in  $U_{1-y}Gd_yO_{2-x}$  solid solutions on Gd content ( $64 \pm 7\text{ cm}^{-1}$  per y) could be used to estimate the Gd content of doped matrices of uncertain composition.

The peak  $\alpha$  at  $\sim 530\text{ cm}^{-1}$  in the broad band in the region  $500$ – $650\text{ cm}^{-1}$  is not observed for pure  $UO_2$  and is attributed to the presence of the oxygen vacancies associated with  $Gd^{3+}$ . Razdan and Shoesmith [25] also observed this peak and attributed it to the creation of oxygen vacancies in Gd-doped and Dy-doped  $UO_2$ , and Talip et al. [27] also demonstrated the existence of oxygen vacancies in  $UO_2$  with 6, 11, 22 mol% La by Raman spectroscopy. These

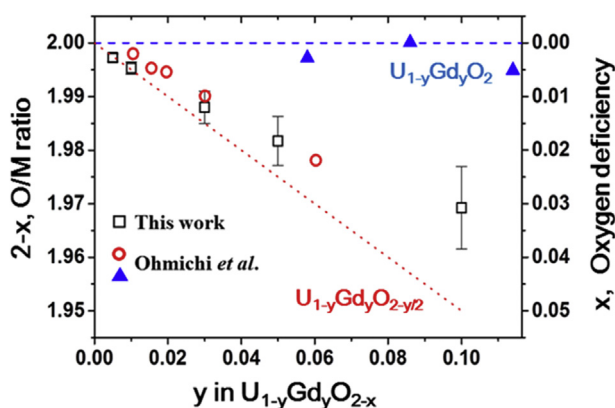
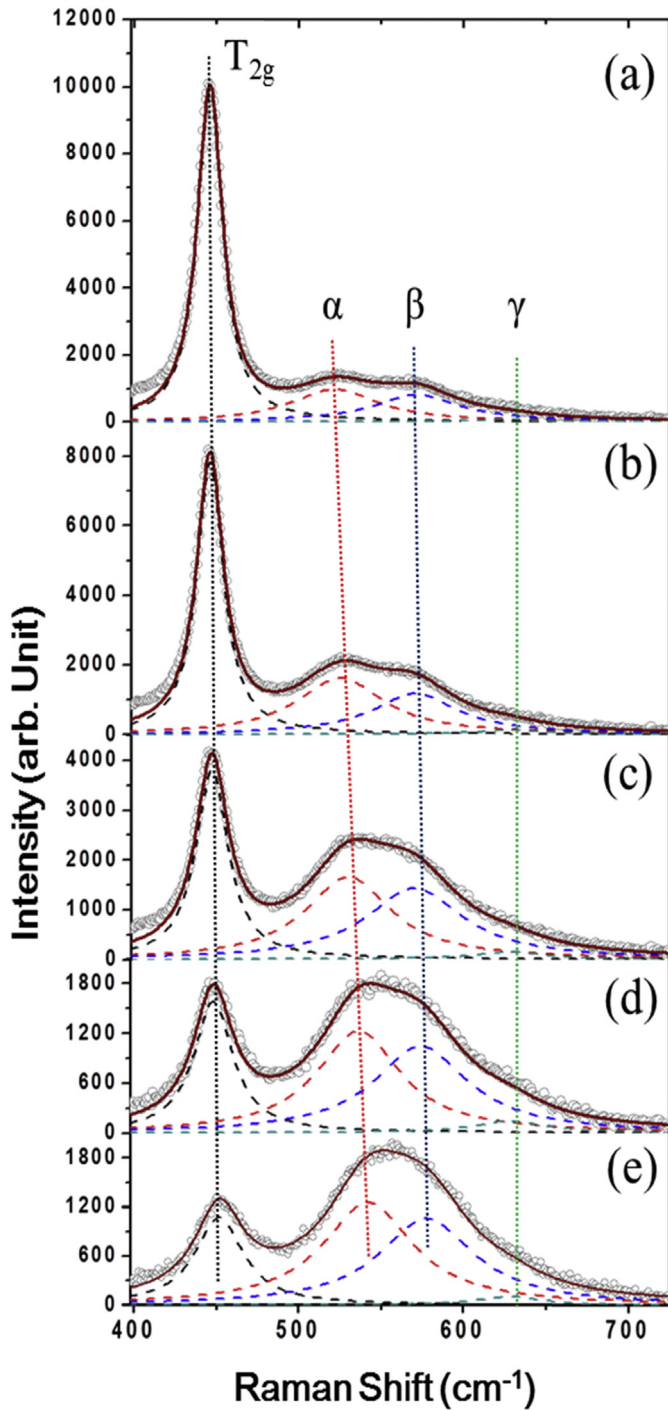
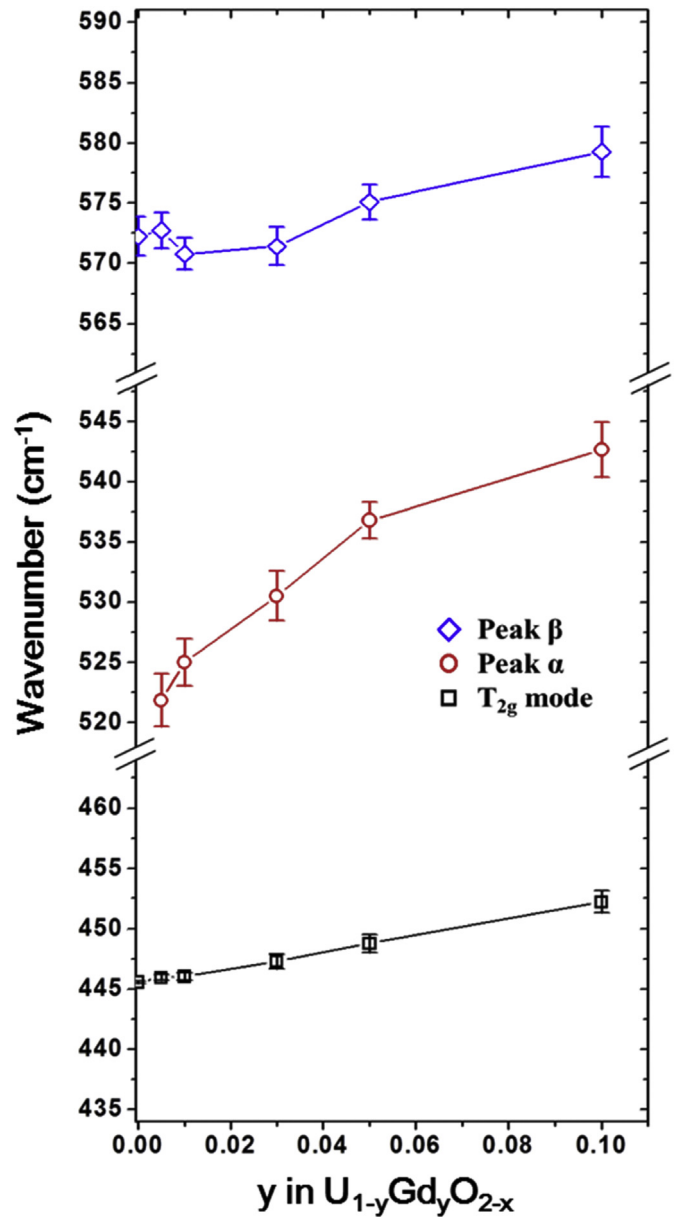


Fig. 2. The black squares show the O/M ratio ( $2-x$ ) and oxygen deficiency ( $x$ ) for  $U_{1-y}Gd_yO_{2-x}$  solid solutions calculated using  $da/dx = 0.024 \pm 0.006$  nm from Ref. [15]. The red circles and blue triangles show the experimental values for  $U_{1-y}Gd_yO_{2-x}$  ( $x \approx y/2$ )-type and  $U_{1-y}Gd_yO_2$  solid solutions, respectively (from Ref. [15]). The red dotted and blue dashed lines show the ideal O/M ratios for perfect  $U_{1-y}Gd_yO_{2-y/2}$  and  $U_{1-y}Gd_yO_2$  solid solutions as a function of Gd content, respectively. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



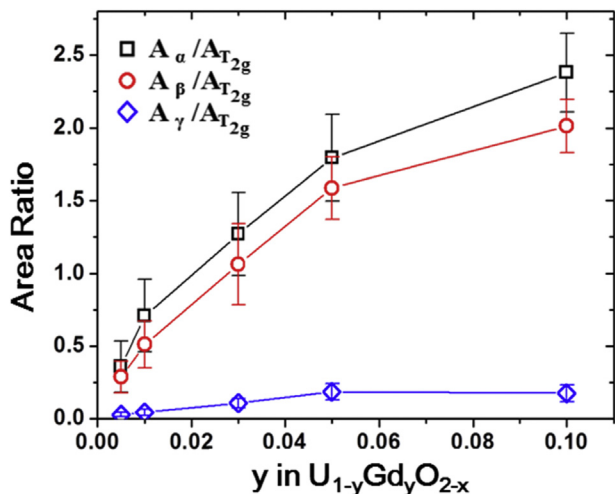
**Fig. 4.** Deconvoluted Raman spectra of (a)  $U_{0.995}Gd_{0.005}O_{2-x}$ , (b)  $U_{0.99}Gd_{0.01}O_{2-x}$ , (c)  $U_{0.97}Gd_{0.03}O_{2-x}$ , (d)  $U_{0.95}Gd_{0.05}O_{2-x}$  and (e)  $U_{0.9}Gd_{0.1}O_{2-x}$  solid solutions with dashed Lorentzian peaks at  $\sim 445$ ,  $530$ ,  $575$  and  $630$   $cm^{-1}$ . Gray dots and brown line represent the experimental data and fitted line, respectively. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

assignments are consistent with those of Desgranges et al. [26] who also interpreted a peak at this wavenumber in spectra measured on  $(U,Nd)O_{2-x}$  as due to a local phonon mode associated with an oxygen-vacancy-induced lattice distortion. A similar peak was observed for  $Ce_{1-y}RE_yO_{2-x}$  samples [36,37] and simulated using a Green's function model calculation [36]. These observations are consistent with the theoretical calculations of Park and Olander [38] for Gd-doped  $UO_{2+x}$  which also indicated the presence of

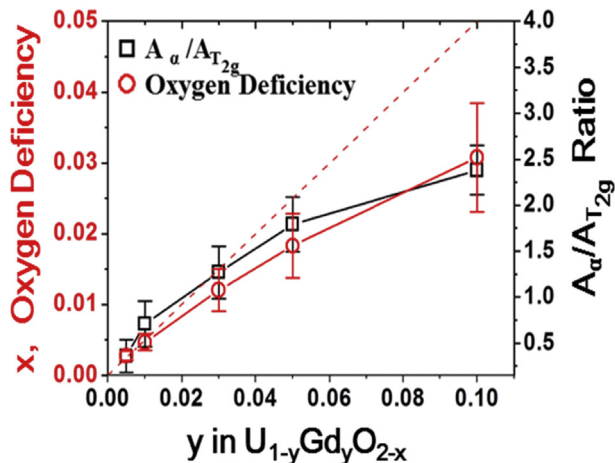


**Fig. 5.** The wavenumber with standard deviation of Lorentzian peak  $T_{2g}$ ,  $\alpha$  and  $\beta$  used to fit Raman spectra of  $U_{1-y}Gd_yO_{2-x}$  solid solutions with Gd contents,  $y = 0.005, 0.01, 0.03, 0.05$  and  $0.1$ .

oxygen vacancies. Additionally, no peak  $\alpha$  was observed in Raman spectra recorded on  $U_{1-y}Th_yO_2$  samples [24], because  $Th^{4+}$  present in solid solution in  $U^{4+}_{1-y}Th^{4+}_yO_{2-x}$  does not require the creation of a significant number of oxygen vacancies to compensate for a charge imbalance. This strongly supports that peak  $\alpha$  is an indication of the presence of oxygen vacancies. In this study, the areas of the  $T_{2g}$ ,  $\alpha$ ,  $\beta$  and  $\gamma$  peaks (denoted as  $A_{T_{2g}}$ ,  $A_\alpha$ ,  $A_\beta$  and  $A_\gamma$ , respectively) were calculated and the area ratios relative to that of the  $T_{2g}$  peak area ( $A_\alpha/A_{T_{2g}}$ ,  $A_\beta/A_{T_{2g}}$  and  $A_\gamma/A_{T_{2g}}$ ) are plotted as a function of Gd content in Fig. 6. The  $A_\alpha/A_{T_{2g}}$  ratio which symbolizes the relative importance of oxygen vacancies in the fluorite lattice structure of  $UO_2$  increases with increasing Gd content. As shown in Fig. 7 this increase is directly related to the increased oxygen deficiency determined using the lattice parameters obtained by XRD to calculate the value of  $x$ , Fig. 2. A similar peak shift was observed for the relative intensity of the  $540$   $cm^{-1}$  band to  $T_{2g}$  band with a  $y$  value of  $U_{1-y}La_yO_{2-y/2}$  [27]. The marked positive shift of peak  $\alpha$



**Fig. 6.** Area ratio between peak  $\alpha$  and  $T_{2g}$  (black square,  $A_\alpha/A_{T_{2g}}$ ), peak  $\beta$  and  $T_{2g}$  (red circle,  $A_\beta/A_{T_{2g}}$ ) and peak  $\gamma$  and  $T_{2g}$  (blue diamond,  $A_\gamma/A_{T_{2g}}$ ) with  $y$  value of  $U_{1-y}Gd_yO_{2-x}$  solid solutions. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



**Fig. 7.** Comparison of the  $A_\alpha/A_{T_{2g}}$  ratio with the calculated oxygen deficiency ( $x$ ) of  $U_{1-y}Gd_yO_{2-x}$  solid solutions with various Gd contents. Red dashed line represents the oxygen deficiency ( $x = y/2$ ) of perfect  $U_{1-y}Gd_yO_{2-y/2}$  solid solutions. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

with increasing Gd content in Fig. 5 is to be expected as a result of the large lattice distortions caused by the large increase in oxygen vacancy concentration or dopant-vacancy clusters, as has been proposed in the mixed oxide defect model [38].

The peak  $\beta$  at  $\sim 575\text{ cm}^{-1}$  has been ascribed to a first order L-O phonon which arises due to crystal lattice disorder [34], although a peak at  $\sim 600\text{ cm}^{-1}$ , observed in a  $Ce_{1-y}RE_yO_{2-x}$  sample [37] has been assigned to the formation of the  $MO_8$ -type complexes not involving oxygen vacancies [37,39,40]. A similar peak observed at  $\sim 575\text{ cm}^{-1}$  for  $UO_2$  samples implanted with Kr and  $He^{2+}$  [32,41,42], and for hyperstoichiometric  $UO_{2+x}$  [23] has been assigned to lattice damage and the disorder induced by the presence of defects, respectively. Thus the increase in the  $A_\beta/A_{T_{2g}}$  ratio with increasing Gd content represents a break down in selection rules due to the existence of oxygen vacancies. This break down of selection rules allows the forbidden first order L-O Raman scattering mode to be observed [43]. The disappearance of the bands in the region from 500 to  $600\text{ cm}^{-1}$  in Raman spectra recorded on stoichiometric

$U_{1-y}Nd_yO_2$  samples [26] and on  $U_{1-y}La_yO_{2-y/2}$  samples during oxidation [27] strongly supports the association of peaks  $\alpha$  and  $\beta$  with lattice defects due to oxygen vacancies. The positive shift of peak  $\beta$  with increasing Gd contents shows a similar tendency to that of peak  $T_{2g}$ .

The peak  $\gamma$  at  $\sim 630\text{ cm}^{-1}$  shown slightly in broad bands of all  $U_{1-y}Gd_yO_{2-x}$  solid solutions has been ascribed to cuboctahedral clusters in the  $U_4O_9$  phase [44]. This assignment was confirmed by Talip et al. during oxidation of  $U_{1-y}La_yO_{2-y/2}$  [27]. Its general absence confirms the absence of this highly oxidized cluster in our specimens. The very slight increase in the  $A_\gamma/A_{T_{2g}}$  ratio as the Gd content increases as shown in Fig. 6 may reflect oxidation during handling or a lack of reducing power during sintering. Such a surface oxidation may also result in the large difference between calculated oxygen deficiencies of  $U_{1-y}Gd_yO_{2-x}$  solid solutions and those of fully reduced  $U_{1-y}Gd_yO_{2-y/2}$  solid solutions (Fig. 2).

#### 4. Conclusions

A series of  $U_{1-y}Gd_yO_{2-x}$  specimens, with  $y = 0.005, 0.01, 0.03, 0.05$  and  $0.1$ , have been analyzed by XRD and Raman spectroscopy techniques.

XRD showed that an increase in Gd content caused a contraction in the fluorite lattice, indicated by the unit cell parameter, and a decrease in the O/M ratio indicating an increase in the oxygen deficiency of the lattice.

The Raman spectra showed an increased distortion of the fluorite lattice as the doping level increased. The shifts in Raman peak positions were consistent with a lattice contraction and with an increasingly defective structure.

An increase in the relative intensity for the Raman peak at  $535\text{ cm}^{-1}$ , attributed to the presence of Gd dopant-oxygen vacancy clusters, compared to that of the peak at  $445\text{ cm}^{-1}$ , attributed to the U-O symmetric stretching mode ( $T_{2g}$ ) of the undisturbed fluorite lattice, confirmed the number of oxygen vacancies increased as the Gd doping level increased.

The position of the  $T_{2g}$  peak was found to vary linearly with Gd content indicating this shift could be used to determine the Gd content.

The characterization of  $(U,Gd)O_{2-x}$  structures provides fundamental data to improve our understanding of spent nuclear fuel and has important implication for the characterization of nuclear fuel materials.

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